

Information Loss Process in NS Equation = Cause of The Chaos. 2013/10/3,22

Long term prediction is our imminent wish, while chaos in NS eqn solution has been intercepting. The cause of chaos is due to **frictional term** $\mu \nabla^2 \mathbf{V}$, which could be showed to become **heat** ($J_i \equiv \mu [(\text{div} \mathbf{V})^2 + (\text{curl} \mathbf{V})^2]$), but not dynamical momentum energy. This is also nothing, but **entropy increasing process** ($dS = dE/T$) which could be a measure for **Chaos (the irreversibility)** in Navier-Stokes Fluid Equation.

[1] : Friction Force as Heat Dissipative Process.

The most simple example of heat dissipating process is dynamics of friction force.

$$\mathbf{F} = d(m\mathbf{V})/dt = -k(\mathbf{V}/V)/V.$$

* \mathbf{V} is velocity vector and \mathbf{F} is a **model** of friction force proportional to $(1/V)$.

$$dE/dt = \langle \mathbf{V} \cdot \mathbf{F} \rangle = d(m\mathbf{V}^2/2)/dt = -k \langle 0 \rangle.$$

$$E(t_1) - E(t_0) = \int_{t_0}^{t_1} dt (dE/dt) = \int_{t_0}^{t_1} dt d(m\mathbf{V}^2/2)/dt = \int_{t_0}^{t_1} dt (-k) = -k(t_1 - t_0) < 0.$$

That is, initial positive kinetic energy ($m\mathbf{V}(t_0)^2/2$) is absorbed as a heat one to stop $m\mathbf{V}$.

Also note **friction force work(energy)** indicates generating **negative energy/time**.

[2] : Energy Theorem in NS-Fluid Equation,

$$(1) D(\rho \mathbf{V})/Dt = \mathbf{f} = \mu \nabla^2 \mathbf{V} - \text{grad} P + 2 \rho_s \boldsymbol{\Omega} \times \mathbf{V} + \rho \mathbf{g}. \quad \text{<NS equation>}$$

Acceleration force = {frictional + pressure + Coriolis + gravity} forces.

(2) **Energy Theorem** : $E(t) - E(0) = \int_0^t dW = \int_0^t dt \mathbf{V} \cdot \mathbf{f}$
 $\int_0^t dt \mathbf{V} \cdot D(\rho \mathbf{V})/Dt = \int_0^t dt \mathbf{V} \cdot \mu \nabla^2 \mathbf{V} - \int_0^t dt \mathbf{V} \cdot \text{grad} P + \int_0^t dt \mathbf{V} \cdot \rho \mathbf{g}.$

☞ : note energy from **Coriolis force** is vanished due to $0 = \langle \mathbf{V} \cdot 2 \rho_s \boldsymbol{\Omega} \times \mathbf{V} \rangle.$
 If $0 = \langle \mathbf{V} \cdot \mathbf{g} \rangle$, also the energy from gravity becomes zero. Ordinal flow \mathbf{V} in ocean and atmosphere is orthogonal to \mathbf{g} , so gravity contribution is **vertical flow** \mathbf{V} only.

$$(a) K \equiv \int_0^t dt \mathbf{V} \cdot [\mathbf{V} (\partial \rho / \partial t + \mathbf{V} \cdot \text{grad} \rho)] + \rho (\partial \mathbf{V} / \partial t + \mathbf{V} \cdot \text{grad} \mathbf{V})] \\ = \int_0^t dt \{ D(\frac{1}{2} \rho \mathbf{V}^2) / Dt + \frac{1}{2} \mathbf{V}^2 (D\rho / Dt) - \frac{1}{2} \rho (D\mathbf{V}^2 / Dt) + \rho \mathbf{V} \cdot (D\mathbf{V} / Dt) \}.$$

$$(b) L \equiv \int_0^t dt \mathbf{V} \cdot \mu \nabla^2 \mathbf{V} = \mu \int_0^t dt [\text{div}(\mathbf{V} \text{div} \mathbf{V} - \mathbf{V} \times \text{curl} \mathbf{V}) - (\text{div} \mathbf{V})^2 - (\text{curl} \mathbf{V})^2] \\ = -\mu \int_0^t dt [(\text{div} \mathbf{V})^2 + (\text{curl} \mathbf{V})^2]$$

$L \equiv$ (vanishing by surface integral)
 - (thermal loss by jet blow) - (thermal loss by fluid stiring).

$$* \nabla^2 \mathbf{V} = \text{grad} \text{div} \mathbf{V} - \text{curl} \text{curl} \mathbf{V}. \quad * \text{div}(\mathbf{A} \times \mathbf{B}) = \mathbf{B} \text{curl} \mathbf{A} - \mathbf{A} \text{curl} \mathbf{B}.$$

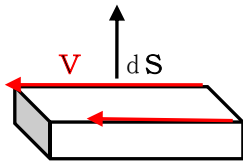
$$-\oint dV \cdot \mathbf{V} \text{curl} \text{curl} \mathbf{V} = -\oint dV \{ \text{div}(\text{curl} \mathbf{V} \times \mathbf{B}) + \text{curl} \mathbf{V} \cdot \text{curl} \mathbf{V} \}.$$

$$\oint dV \cdot \mathbf{V} \text{grad} \text{div} \mathbf{V} = \oint dV \{ \text{div}(\mathbf{V} \text{div} \mathbf{V}) - \text{div} \mathbf{V} \cdot \text{div} \mathbf{V} \}.$$

Thus energy from **friction term** become negative, which mean **energy absorbing** from left momentum energy and the actuality is **dissipative heat by frictional turbulence**.

Note this is just the origin of **chaos** in NS equation !!.

$$\iiint_V dV \operatorname{div}(\mathbf{V} \operatorname{div} \mathbf{V} - \mathbf{V} \times \operatorname{curl} \mathbf{V}) = \iint_S dS (\mathbf{V} \operatorname{div} \mathbf{V} - \mathbf{V} \times \operatorname{curl} \mathbf{V}) = 0.$$



Boundary surface flow \mathbf{V} is orthogonal to dS , so 1st term is zero. $\operatorname{curl} \mathbf{V}$ is parallel with dS , and $\mathbf{V} \times \operatorname{curl} \mathbf{V}$ is orthogonal to dS , so 2nd term is also zero.

$$(c) M \equiv - \int_0^t dt \mathbf{V} \cdot \operatorname{grad} P = - \int_0^t dt [\operatorname{div}(\mathbf{V} \cdot P) - P \operatorname{div} \mathbf{V}].$$

$M \equiv$ (1st term is net work by pressure on a closed surface integral and 2nd one is volume integral energy jet flow ($\operatorname{div} \mathbf{V}$) with pressure).

*What causes pressure ? in ocean and atmosphere flows,

$$(d) N \equiv \int_0^t dt \mathbf{V} \cdot \rho \mathbf{g}. \text{ <This term could be potential energy of } \rho g h, \text{ where } h \text{ is height (depth)>.$$

This is a work by **gravity** with parallel component of \mathbf{V} . This is a vertical flow due to **sink** (ρ change by dense **salinity** in sea ice freezing), or ρ change by **cold heat (convection)** or **buoyancy** force (density ρ change by **hot heat (convection)**).

[3] : Entropy Generating in NS-Fluid Equation

of unknown variables $\{V, P, \rho, T, Q\}$. <independent eqns {NS, (2)(3)(4)(5)}>

-The simultaneous equations-

$$E(t) - E(0) = \int_0^t dW = \int_0^t dt \mathbf{V} \cdot \mathbf{f}.$$

Fluid dynamical energy relation $\equiv W$:

$$\begin{aligned} dW/dt &= \mathbf{V} \cdot D(\rho \mathbf{V})/Dt = \mathbf{V} \cdot \mu \nabla^2 \mathbf{V} - \mathbf{V} \cdot \operatorname{grad} P + \mathbf{V} \cdot \rho \mathbf{g}. \\ &= -\mu [(\operatorname{div} \mathbf{V})^2 + (\operatorname{curl} \mathbf{V})^2] - \mathbf{V} \cdot \operatorname{grad} P + \mathbf{V} \cdot \rho \mathbf{g}. \end{aligned}$$

Thermal energy relation $\equiv Q$:

$$\begin{aligned} \partial Q(\mathbf{x}, t) / \partial t &= -\operatorname{div}(\mathbf{Q} \mathbf{V}) + J_i + \lambda(\mathbf{x}, t) \operatorname{grad} T(\mathbf{x}, t) + J_R(\mathbf{x}, t). \\ J_i &\equiv +\mu [(\operatorname{div} \mathbf{V})^2 + (\operatorname{curl} \mathbf{V})^2]. \end{aligned}$$

Total Energy Relation:Energy Conservation Low in NS Fluid-Equation $E \equiv W + Q$.

$$(d/dt) [W + Q] = -\mathbf{V} \cdot \text{grad} P + \mathbf{V} \cdot \rho \mathbf{g} - \text{div}(\mathbf{Q} \mathbf{V}) + \lambda(\mathbf{x}, t) \text{grad} T(\mathbf{x}, t) + J_R(\mathbf{x}, t).$$

system energy = {pressure + gravity potential} + {convection + conduction + radiation} heats

(1) **2nd low of Thermodynamics <entropy density increasing>**:

$$dS = dQ/T. \quad dS/dt = (dQ/dt)/T.$$

(2) **Thermodynamical state equation of fluid(ideal gas):**

$$P V = nRT. \quad R \equiv \text{gas constant. } M = \text{molecular mass.}$$

$$P = (m/V M)RT = (\rho / M)RT.$$

(3) **Specific heat C_P of fluid matter.**

$$Q(\mathbf{x}, t) = C_P T(\mathbf{x}, t).$$

(4) **Heat flow conservation low:**

$$\begin{aligned} \partial Q(\mathbf{x}, t) / \partial t &= J_i \{ \text{internal heat generation by } \mu \nabla^2 \mathbf{V} \} + \\ &+ J \{ \text{heat flows by external radiation, conduction and convection} \} \\ &= J_i - \text{div}(\mathbf{Q} \mathbf{V}) + J \{ \text{heat input by external radiation, conduction} \}. \end{aligned}$$

(a) **radiation** $\equiv J_R(\mathbf{x}, t)$: author don't know the explicit form !

(b) **conduction** $\equiv J_{cd}(\mathbf{x}, t)$: $\lambda(\mathbf{x}, t) \text{grad} T(\mathbf{x}, t)$

(c) **convection** $\equiv J_{cv}(\mathbf{x}, t)$: $-\text{div}(\mathbf{Q}(\mathbf{x}, t) \mathbf{V}(\mathbf{x}, t))$

(5) **fluid mass(ocean or atmosphere)density ρ conservation low:**

$$\partial \rho(\mathbf{x}, t) / \partial t = -\text{div}(\rho \mathbf{V}) + J_m \{ ? \}.$$

Those simultaneous equations {NS, (2)(3)(4)(5)} could determine unknown variables {V, P, ρ , T, Q}. Thus we could calculate (dS/dt) **the entropy increasing.**

[4] : **Information measure in NS-Fluid Equation.**

Reference: Le'on Brillouin, Science and Information Theory, Academic Press, N. Y., 1956, 1962.

In this book, Brillouin described "man observing information" is measured by negative entropy. Therefore entropy increasing in NS Fluid Equation is not reversible, but **irreversible generating positive entropy** which means nothing but "information loss" in Fluid equation solution (**Chaos**). Therefore long time prediction could not be accomplished without reducing friction term $\mu \nabla^2 \mathbf{V}$, which could be done by so called **scale transformation in time and space variables.**

APPENDIX_1:

How to transform Fluid Equation more causalitical !! 2013-3-27,9-26.

Large space and time scale view by scaling transformation(Reinolds analogy low).

So called Chaos in fluid equation is due to friction term $\mu \nabla^2 \mathbf{u}$, which causes long term weather prediction difficult. However a method could overcome the difficulty.

Fluid equation could be transformed into new space and time variables $\{x', t'\}$ from $\{x, t\}$ by $\{x \equiv Lx', t \equiv (L/U)t', \rho \equiv \rho'; m \equiv L^3 m', \Omega \equiv \Omega'\}$

$$\begin{aligned} u &\equiv \partial x / \partial t = \partial (Lx') / \partial ((L/U)t') = u \partial x' / \partial t' = Uu' . \\ \alpha &\equiv \partial u / \partial t = \partial (Uu') / \partial ((L/U)t') = (U^2/L) \partial u' / \partial t' = (U^2/L) \alpha' . \\ P &\equiv f/x^2 = m \alpha / x^2 = L^3 m' (U^2/L) \alpha' / (Lx')^2 = U^2 m' \alpha' / x'^2 = U^2 p' . \\ \partial P / \partial x &= (U^2/L) \partial p' / \partial x' . \\ \rho &\equiv m/x^3 = L^3 m' / (Lx')^3 = m' / x'^3 = \rho' . \\ \partial^2 u / \partial x^2 &= \partial^2 (Uu') / \partial (Lx')^2 = (U/L^2) \partial^2 u' / \partial x'^2 \\ -2 \rho \Omega \times \mathbf{u} &= -(U) 2 \rho' \Omega' \times \mathbf{u}' \end{aligned}$$

$$\begin{aligned} \rho \{ \partial \mathbf{u} / \partial t + \mathbf{u} \cdot \text{grad} \mathbf{u} \} &= -\text{grad} P + \mu \nabla^2 \mathbf{u} - 2 \rho \Omega \times \mathbf{u} + \rho \mathbf{g} . \\ \rho' \{ (U^2/L) \partial \mathbf{u}' / \partial t' + (U^2/L) \mathbf{u}' \cdot \text{grad}' \mathbf{u}' \} & \\ = -(U^2/L) \text{grad}' P' + \mu (U/L^2) \nabla'^2 \mathbf{u}' - (U) 2 \rho' \Omega' \times \mathbf{u}' + \rho' \mathbf{g} & . \end{aligned}$$

$$\begin{aligned} \rho' \{ \partial \mathbf{u}' / \partial t' + \mathbf{u}' \cdot \text{grad}' \mathbf{u}' \} & \\ = -\text{grad}' P' + (\mu / UL) \nabla'^2 \mathbf{u}' - 2 \rho' (L/U) \Omega' \times \mathbf{u}' + \rho' (L/U^2) \mathbf{g} & \\ = -\text{grad}' (P/U^2) + (\mu / UL) \nabla'^2 \mathbf{u}' - 2 \rho' (L/U) \Omega' \times \mathbf{u}' + \rho' (L/U^2) \mathbf{g} . & \end{aligned}$$

Then taking $U=1, L=\text{larger}$ could make $(\mu / UL) \equiv \mu'$ smaller(almost nothing friction force).
grad' u', Pressure gradient, Coliori force and gravity become larger.

This is to make fluid equation more causalitical one in **large space and time scale view**.

This might be a certification for global climate model calculation.

example) :

$$\begin{aligned} 60\text{km (man size)} &\rightarrow 6 \times 10^6 \text{m (earth size)} \Rightarrow U \equiv 1; L \equiv 10^2 . \\ \mu = 1.307 \times 10^{-3} \text{Ns/m}^2 &\rightarrow \mu' = \mu / UL = 1.307 \times 10^{-5} \text{Ns/m}^2 . \end{aligned}$$

$$\rho' \{ \partial \mathbf{u}' / \partial t' + \mathbf{u}' \cdot \text{grad}' \mathbf{u}' \} = -\text{grad}' (P') - 2 \rho' (L/U) \Omega' \times \mathbf{u}' + \rho' (L/U^2) \mathbf{g} .$$

In order to be **complete dynamic equation**, left term need $+u \{ \partial \rho / \partial t + \mathbf{u} \cdot \text{grad} \rho \}$ for force by density ρ' change.

Appendix _2 : :vector calculus in left term.

$$\begin{aligned}
 D(\rho \mathbf{V})/Dt &= \partial_t(\rho \mathbf{V}) + \langle \mathbf{V} \cdot \text{grad}(\rho \mathbf{V}) \rangle \\
 &= \partial_t(\rho \mathbf{V}) + \langle \mathbf{V} \cdot (\text{grad}(\rho \mathbf{V})) \rangle = \partial_t(\rho \mathbf{V}) + \mathbf{V} \langle \mathbf{V} \cdot \text{grad} \rho \rangle + \rho \langle \mathbf{V} \cdot \text{grad} \mathbf{V} \rangle \\
 &= \underline{\rho [\partial_t \mathbf{V} + \langle \mathbf{V} \cdot \text{grad} \mathbf{V} \rangle]} + \mathbf{V} [\partial_t \rho + \langle \mathbf{V} \cdot \text{grad} \rho \rangle] \dots \dots \text{OK}
 \end{aligned}$$

$$\begin{aligned}
 * (\text{grad}(\rho \mathbf{V})) &\equiv [\text{grad}(\rho V_1), \text{grad}(\rho V_2), \text{grad}(\rho V_3)] \\
 &= [V_1 \text{grad} \rho + \rho \text{grad} V_1, V_2 \text{grad} \rho + \rho \text{grad} V_2, V_3 \text{grad} \rho + \rho \text{grad} V_3] \\
 &= [V_1 \text{grad} \rho + \rho \text{grad} V_1, V_2 \text{grad} \rho + \rho \text{grad} V_2, V_3 \text{grad} \rho + \rho \text{grad} V_3] \\
 \mathbf{V} \cdot (\text{grad}(\rho \mathbf{V})) &= \mathbf{V} \langle \mathbf{V} \cdot \text{grad} \rho \rangle + \rho \langle \mathbf{V} \cdot \text{grad} \mathbf{V} \rangle. \dots \dots \text{OK}
 \end{aligned}$$

$$\begin{aligned}
 * * \underline{\rho \mathbf{V} \cdot (\mathbf{V} \cdot \text{grad} \mathbf{V})} &= \langle \mathbf{V} \cdot \langle \mathbf{V} \cdot (\text{grad}(\rho \mathbf{V})) \rangle \rangle - \mathbf{V}^2 \langle \mathbf{V} \cdot \text{grad} \rho \rangle \\
 \mathbf{K} &\equiv \mathbf{V} \cdot [\mathbf{V} (\partial_t \rho / \partial t + (\mathbf{V} \cdot \text{grad} \rho)) + \rho (\partial_t \mathbf{V} / \partial t + \mathbf{V} \cdot \text{grad} \mathbf{V})] \\
 &= \mathbf{V}^2 (\partial_t \rho / \partial t) + (\mathbf{V} \cdot \text{grad} \rho) + \frac{1}{2} \rho \partial_t (\mathbf{V}^2) + \rho \mathbf{V} \cdot (\mathbf{V} \cdot \text{grad} \mathbf{V}) \\
 &= \partial_t (\frac{1}{2} \rho \mathbf{V}^2) + \frac{1}{2} \mathbf{V}^2 \partial_t \rho + \mathbf{V}^2 (\mathbf{V} \cdot \text{grad} \rho) + \rho \mathbf{V} \cdot (\mathbf{V} \cdot \text{grad} \mathbf{V}) \\
 &= \underline{\partial_t (\frac{1}{2} \rho \mathbf{V}^2) + (\frac{1}{2} \mathbf{V}^2) \langle \mathbf{V} \cdot \text{grad} \rho \rangle + \rho \mathbf{V}^2 \text{div} \mathbf{V}} \\
 &+ \frac{1}{2} \mathbf{V}^2 \partial_t \rho + \mathbf{V}^2 (\mathbf{V} \cdot \text{grad} \rho) + \rho \mathbf{V} \cdot (\mathbf{V} \cdot \text{grad} \mathbf{V}) \\
 &- (\frac{1}{2} \mathbf{V}^2) \langle \mathbf{V} \cdot \text{grad} \rho \rangle - \rho \mathbf{V}^2 \text{div} \mathbf{V} \\
 &= D((\frac{1}{2} \rho \mathbf{V}^2)/Dt + \frac{1}{2} \mathbf{V}^2 \partial_t \rho + (\frac{1}{2} \mathbf{V}^2) \langle \mathbf{V} \cdot \text{grad} \rho \rangle + \rho \mathbf{V} \cdot (\mathbf{V} \cdot \text{grad} \mathbf{V}) - \rho \mathbf{V}^2 \text{div} \mathbf{V} \\
 &= D((\frac{1}{2} \rho \mathbf{V}^2)/Dt + \frac{1}{2} \mathbf{V}^2 D \rho / Dt + \underline{\rho \mathbf{V} \cdot [(\mathbf{V} \cdot \text{grad} \mathbf{V}) - \mathbf{V} \text{div} \mathbf{V}]}
 \end{aligned}$$

$$\begin{aligned}
 D((\frac{1}{2} \rho \mathbf{V}^2)/Dt &= \partial_t (\frac{1}{2} \rho \mathbf{V}^2) + \langle \mathbf{V} \cdot \text{grad}(\frac{1}{2} \rho \mathbf{V}^2) \rangle \\
 &= \partial_t (\frac{1}{2} \rho \mathbf{V}^2) + \langle \mathbf{V} \cdot [(\frac{1}{2} \mathbf{V}^2 \cdot \text{grad} \rho) + \rho \mathbf{V} \text{div} \mathbf{V}] \rangle \\
 &= \partial_t (\frac{1}{2} \rho \mathbf{V}^2) + (\frac{1}{2} \mathbf{V}^2) \langle \mathbf{V} \cdot \text{grad} \rho \rangle + \rho \mathbf{V}^2 \text{div} \mathbf{V}. \dots \dots \text{OK}
 \end{aligned}$$

$$\text{grad} \mathbf{V}^2 = 2 \mathbf{V} \text{div} \mathbf{V}.$$

$$\begin{aligned}
 D\mathbf{V}^2/Dt &= \partial_t \mathbf{V}^2 + \langle \mathbf{V} \cdot \text{grad} \mathbf{V}^2 \rangle = \partial_t \mathbf{V}^2 + \langle \mathbf{V} \cdot 2 \mathbf{V} \text{div} \mathbf{V} \rangle = \partial_t \mathbf{V}^2 + 2 \mathbf{V}^2 \text{div} \mathbf{V}. \\
 &= 2 \mathbf{V} \partial_t \mathbf{V} + 2 \mathbf{V}^2 \text{div} \mathbf{V} = 2 \mathbf{V} (\partial_t \mathbf{V} + \mathbf{V} \text{div} \mathbf{V}). \\
 \frac{1}{2} D\mathbf{V}^2/Dt &= \mathbf{V} (\partial_t \mathbf{V} + \mathbf{V} \text{div} \mathbf{V}). \\
 \frac{1}{2} \rho D\mathbf{V}^2/Dt &= \rho \mathbf{V} (\partial_t \mathbf{V} + \mathbf{V} \text{div} \mathbf{V}).
 \end{aligned}$$

$$-\rho \mathbf{V} \cdot (\mathbf{V} \operatorname{div} \mathbf{V}) = \rho \mathbf{V} \cdot \partial_t \mathbf{V} - \frac{1}{2} \rho D \mathbf{V}^2 / Dt$$

$$(\mathbf{V} \cdot \operatorname{grad} \mathbf{V}) = D \mathbf{V} / Dt - \partial_t \mathbf{V}$$

$$\rho \mathbf{V} \cdot (\mathbf{V} \cdot \operatorname{grad} \mathbf{V}) = \rho \mathbf{V} \cdot (D \mathbf{V} / Dt - \partial_t \mathbf{V})$$

$$\rho \mathbf{V} \cdot [(\mathbf{V} \cdot \operatorname{grad} \mathbf{V}) - \mathbf{V} \operatorname{div} \mathbf{V}] = -\frac{1}{2} \rho D \mathbf{V}^2 / Dt + \rho \mathbf{V} \cdot (D \mathbf{V} / Dt$$

$$\begin{aligned} K &= D \left(\frac{1}{2} \rho \mathbf{V}^2 \right) / Dt + \frac{1}{2} \mathbf{V}^2 D \rho / Dt + \rho \mathbf{V} \cdot [(\mathbf{V} \cdot \operatorname{grad} \mathbf{V}) - \mathbf{V} \operatorname{div} \mathbf{V}] \\ &= D \left(\frac{1}{2} \rho \mathbf{V}^2 \right) / Dt + \frac{1}{2} \mathbf{V}^2 (D \rho / Dt) - \frac{1}{2} \rho (D \mathbf{V}^2 / Dt) + \rho \mathbf{V} \cdot (D \mathbf{V} / Dt). \end{aligned}$$

* Now author don't know well about the physical meanings.